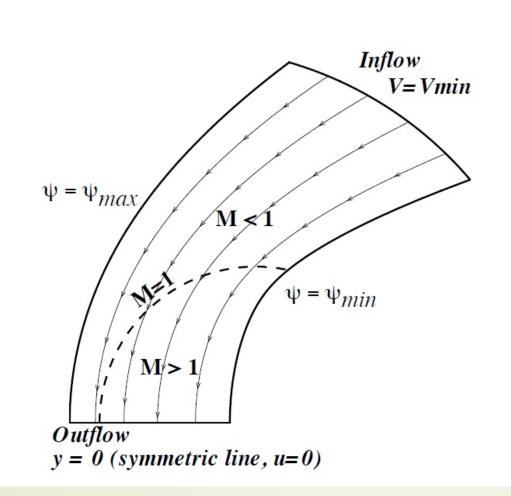
### Shallow Water Equations Ringleb's Flow

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#### Ringleb's flow

- A particular flow where the fluid accelerates from subsonic to supersonic regime smoothly Without any shock.
- Much Useful to test the accuracy of an inviscid flow for both subsonic and supersonic flow as the exact solution is available.
- Limited to gas flow with  $\gamma = 1.4$ , governed by Euler equations
- A good candidate for order of accuracy test on shallow water equations.

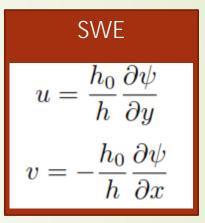
#### Ringleb's flow



Ringleb Flow: A flow in a channel, being smoothly accelerated from subsonic to supersonic. The supersonic region can be adjusted by varying  $\psi$ 

Start with potential and stream equations,

$$u = \frac{\partial \phi}{\partial x} \qquad \qquad u = \frac{\rho_0}{\rho} \frac{\partial \psi}{\partial y}$$
$$v = \frac{\partial \phi}{\partial y} \qquad \qquad v = -\frac{\rho_0}{\rho} \frac{\partial \psi}{\partial x}$$



• Introduce flow quantities, V and  $\theta$ 

$$u = V cos\theta$$
$$v = V sin\theta$$

$$\begin{aligned} d\phi &= \phi_x dx + \phi_y dy = V(\cos\theta dx + \sin\theta dy) \\ d\psi &= \psi_x dx + \psi_y dy = \left(\frac{\rho}{\rho_0}\right) V(-\sin\theta dx + \cos\theta dy) \\ dx &= \frac{\cos\theta}{V} d\phi - \frac{\rho_0}{\rho} \frac{\sin\theta}{V} d\psi \\ dy &= \frac{\sin\theta}{V} d\phi + \frac{\rho_0}{\rho} \frac{\cos\theta}{V} d\psi \end{aligned}$$

 $\phi(x,y),\psi(x,y)$ 

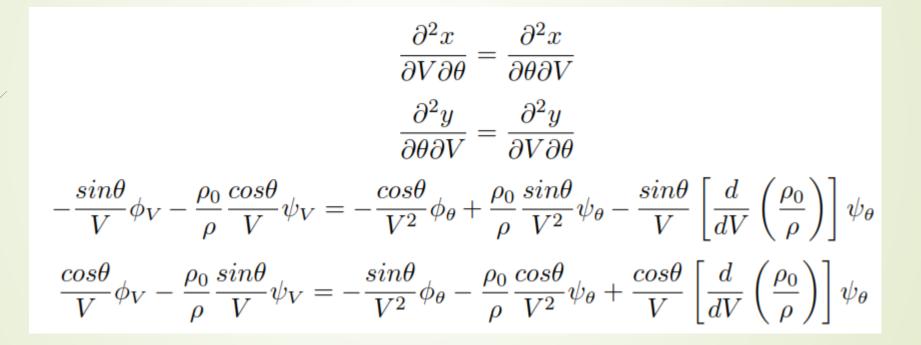
$$d\phi = \frac{\partial \phi}{\partial V} dV + \frac{\partial \phi}{\partial \theta} d\theta = \phi_V dV + \phi_\theta d\theta$$
$$d\psi = \frac{\partial \psi}{\partial V} dV + \frac{\partial \psi}{\partial \theta} d\theta = \psi_V dV + \psi_\theta d\theta$$

$$dx = \left(\frac{\cos\theta}{V}\phi_V - \frac{\rho_0}{\rho}\frac{\sin\theta}{V}\psi_V\right)dV + \left(\frac{\cos\theta}{V}\phi_\theta - \frac{\rho_0}{\rho}\frac{\sin\theta}{V}\psi_\theta\right)d\theta$$

$$dy = (\frac{\sin\theta}{V}\phi_V + \frac{\rho_0}{\rho}\frac{\cos\theta}{V}\psi_V)dV + (\frac{\sin\theta}{V}\phi_\theta + \frac{\rho_0}{\rho}\frac{\cos\theta}{V}\psi_\theta)d\theta$$

V and  $\theta$  are independent of each other.

$$\begin{aligned} \frac{\partial x}{\partial V} &= x_V = \frac{\cos\theta}{V}\phi_V - \frac{\rho_0}{\rho}\frac{\sin\theta}{V}\psi_V\\ \frac{\partial x}{\partial \theta} &= x_\theta = \frac{\cos\theta}{V}\phi_\theta - \frac{\rho_0}{\rho}\frac{\sin\theta}{V}\psi_\theta\\ \frac{\partial y}{\partial V} &= y_V = \frac{\sin\theta}{V}\phi_V + \frac{\rho_0}{\rho}\frac{\cos\theta}{V}\psi_V\\ \frac{\partial y}{\partial \theta} &= y_\theta = \frac{\sin\theta}{V}\phi_\theta + \frac{\rho_0}{\rho}\frac{\cos\theta}{V}\psi_\theta\end{aligned}$$



multiply by  $cos\theta$  ,  $sin\theta$  then sum,

$$\phi_{\theta} = \left(\frac{\rho_0}{\rho}\right) V \psi_V$$

multiply by  $sin\theta$ ,  $cos\theta$  then subtract,

$$\phi_V = \left[ -\frac{\rho_0}{\rho} \frac{1}{V} + \frac{d}{dV} \left( \frac{\rho_0}{\rho} \right) \right] \psi_\theta = V \left[ \frac{d}{dV} \left( \frac{1}{V} \frac{\rho_0}{\rho} \right) \right] \psi_\theta$$

For isentropic flow,  $\frac{dp}{d\rho} = c^2$  (speed of sound) Euler's equation for frictionless, irrotational motion,

 $dp = -\rho V dV$  $\frac{d}{dV} \left(\frac{\rho_0}{\rho}\right) = -\frac{\rho_0}{\rho^2} \frac{d\rho}{dV} = -\frac{\rho_0}{\rho^2} \frac{d\rho}{dp} \frac{dp}{dV}$  $= -\frac{\rho_0}{\rho^2} \left(\frac{1}{c^2}\right) (-\rho V) = -\frac{\rho_0}{\rho^2} \frac{V}{c^2}$ 

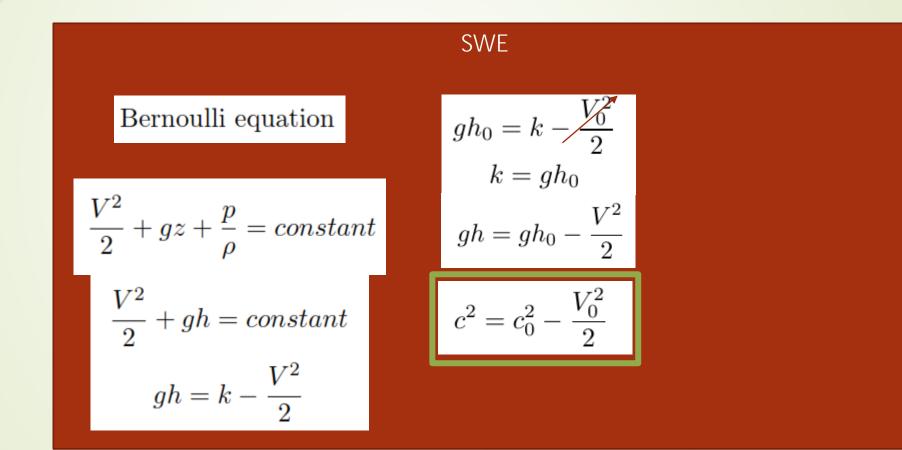
SWE  
Bernoulli equation
$$\begin{aligned}
h &= k - \frac{V^2}{2g} \\
\frac{V^2}{2} + gz + \frac{p}{\rho} = constant \\
\frac{V^2}{2} + gh = constant \\
gh &= k - \frac{V^2}{2} \\
\end{aligned}$$

$$\begin{aligned}
h &= k - \frac{V^2}{2g} \\
\frac{dh}{dV} &= -\frac{1}{g}V \\
\frac{d}{dV} \left(\frac{h_0}{h}\right) &= -\frac{h_0}{h^2}\frac{dh}{dV} \\
&= \frac{h_0}{h^2} \left(-\frac{V}{g}\right) &= \frac{h_0}{h}\frac{V}{gh} \\
\end{aligned}$$

$$c &= \sqrt{gh}$$

$$\begin{split} \phi_V &= -\frac{\rho_0}{\rho} \frac{1}{V} \left( 1 - \frac{V^2}{c^2} \right) \psi_\theta \\ \phi_{\theta V} &= \phi_{V\theta} \\ V^2 \psi_{VV} + V \left( 1 + \frac{V^2}{c^2} \right) \psi_V + \left( 1 - \frac{V^2}{c^2} \right) \psi_{\theta\theta} = 0 \\ c^2 &= c_0^2 - \frac{k - 1}{2} V^2 \end{split}$$

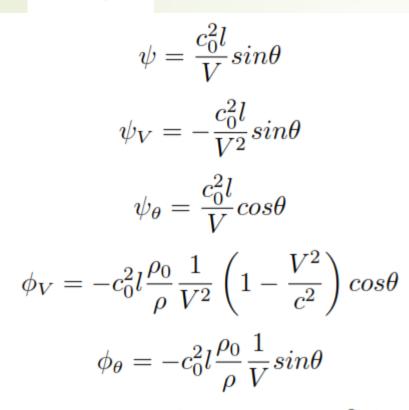
$$V^{2}\left(1-\frac{k-1}{2}\frac{V^{2}}{c_{0}^{2}}\right)\psi_{VV} + V\left(1-\frac{k-3}{2}\frac{V^{2}}{c_{0}^{2}}\right)\psi_{V} + \left(1-\frac{k+1}{2}\frac{V^{2}}{c_{0}^{2}}\right)\psi_{\theta\theta} = 0$$
(1)



$$\begin{split} \psi &= P(V) \cdot Q(\theta) \\ \psi_V &= QP'; \psi_{VV} = QP"; \psi_{\theta\theta} = PQ" \\ V^2 \left[ \frac{1 - \frac{k-1}{2} \frac{V^2}{c_0^2}}{1 - \frac{k+1}{2} \frac{V^2}{c_0^2}} \right] \frac{P"}{P} + V \left[ \frac{1 - \frac{k-3}{2} \frac{V^2}{c_0^2}}{1 - \frac{k+1}{2} \frac{V^2}{c_0^2}} \right] \frac{P'}{P} + \frac{Q"}{Q} = 0 \end{split}$$

V and  $\theta$  independent ;  $\frac{Q^{"}}{Q} = -n^2$  (simple harmonic equation)  $n=1, P = \frac{1}{V}$ 

Hence, let



$$\begin{aligned} x_V &= c_0^2 l \frac{\rho_0}{\rho} \left( -\frac{\cos 2\theta}{V^2} + \frac{\cos^2 \theta}{Vc^2} \right) \\ x_\theta &= -c_0^2 l \frac{\rho_0}{\rho} \frac{\sin 2\theta}{V^2} \\ y_V &= c_0^2 l \frac{\rho_0}{\rho} \left( -\frac{\sin 2\theta}{V^2} + \frac{\sin \theta \cos \theta}{Vc^2} \right) \\ y_\theta &= c_0^2 l \frac{\rho_0}{\rho} \frac{\cos 2\theta}{V^2} \end{aligned}$$

$$x = \int x_{\theta} d\theta + f(V) = c_0^2 l \frac{\rho_0}{\rho} \frac{\cos 2\theta}{V^2} + f(V)$$

$$x_{V} = \frac{c_{0}^{2}l}{2}cos2\theta \left[ -\frac{\rho_{0}}{\rho}\frac{2}{V^{3}} + \frac{1}{V^{2}}\left(\frac{\rho_{0}}{\rho}\frac{V}{c^{2}}\right) \right] + f'(V)$$
$$f'(V) = \frac{c_{0}^{2}}{2}\frac{\rho_{0}}{\rho}\frac{1}{Vc^{2}}$$
$$f(V) = \frac{c_{0}^{2}}{2}\int\frac{\rho_{0}}{\rho}\frac{dV}{Vc^{2}} = \frac{l}{2}L$$
$$L = -\left[\frac{1}{2}\ln\frac{1+\gamma}{1-\gamma} - \frac{1}{\gamma} - \frac{1}{3\gamma^{2}} - \frac{1}{5\gamma^{5}}\right]$$

SWE  

$$f'(V) = \frac{c_0^2}{2} \frac{h_0}{h} \frac{1}{Vc^2}$$

$$f(V) = \frac{c_0^2}{2} \int \frac{h_0}{h} \frac{dV}{Vc^2}$$

$$= \frac{1}{2} \int \frac{h_0}{h} \frac{gh_0}{gh} \frac{dV}{V} = \frac{l}{2}L$$

$$L = \frac{1}{2} \left[ \frac{1}{b} + \ln \frac{1-b}{b} \right]$$

$$b = \frac{h}{h_0} = 1 - \frac{V^2}{2c_0^2}$$

$$\begin{split} \frac{x}{l} &= \frac{1}{2} \left[ \frac{\rho_0}{\rho} \left( \frac{c_0}{V} \right)^2 \cos\theta + L \right] \\ \frac{y}{l} &= \frac{1}{2} \frac{\rho_0}{\rho} \left( \frac{c_0}{V} \right)^2 \sin 2\theta \\ \frac{x}{l} &= \frac{1}{2} \frac{\rho_0}{\rho} \left( \frac{c_0^2}{V^2} - 2\frac{\psi^2}{l^2 c_0^2} \right) + \frac{L}{2} \\ \frac{y}{l} &= \pm \frac{\rho_0}{\rho} \frac{c_0}{V} \frac{\psi}{c_0 l} \sqrt{1 - \frac{V^2}{c_0^2} \left( \frac{\psi}{l c_0} \right)^2} \end{split}$$

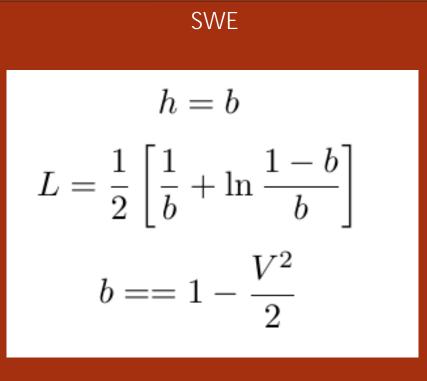
Equations to build up Ringleb's flow domain

The following equations are required to be implemented in Euler code to compute positions of boundary and exact solution.

$$\begin{split} \psi &= \frac{\sin \theta}{V}, \\ x(\psi, V) &= \frac{1}{\rho} \left[ \frac{1}{2V^2} - \psi^2 \right] + \frac{L}{2}, \\ y(\psi, V) &= \pm \frac{\psi}{\rho V} \sqrt{1 - V^2 \psi^2}, \\ \end{split} \qquad \begin{array}{l} \rho &= b^5, \\ L &= \frac{1}{b} + \frac{1}{3b^3} + \frac{1}{5b^5} - \frac{1}{2} \ln \left( \frac{1 + b}{1 - b} \right), \\ b &= \sqrt{1 - 0.2V^2}, \\ p &= b^7. \end{split}$$

$$\begin{split} \rho &= b^5, \\ L &= \frac{1}{b} + \frac{1}{3b^3} + \frac{1}{5b^5} - \frac{1}{2} \ln \left( \frac{1+b}{1-b} \right), \\ b &= \sqrt{1-0.2V^2}, \end{split}$$

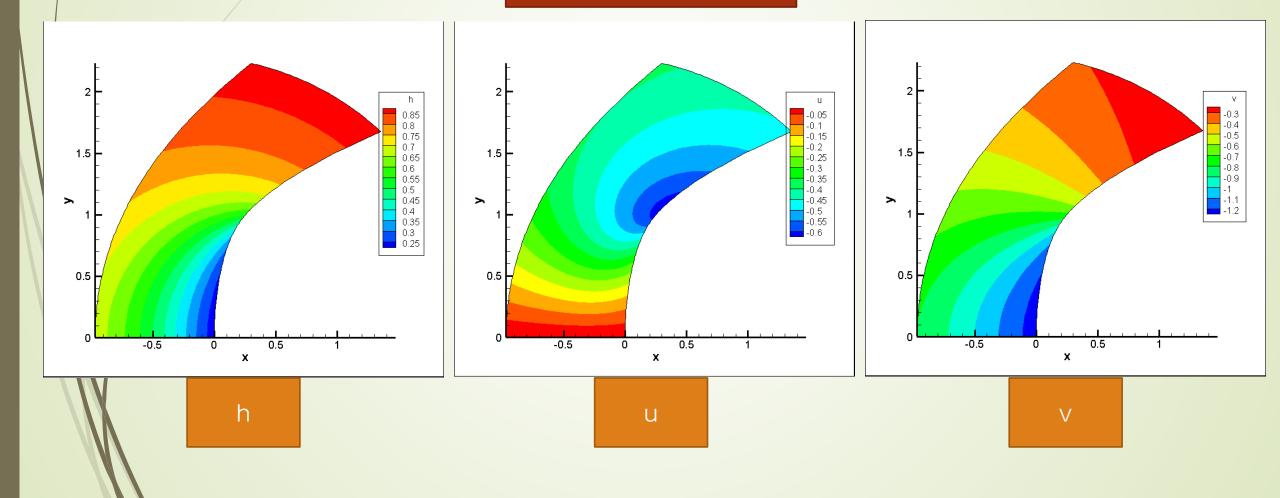
Variables have been non-dimensionalized by stagnation values



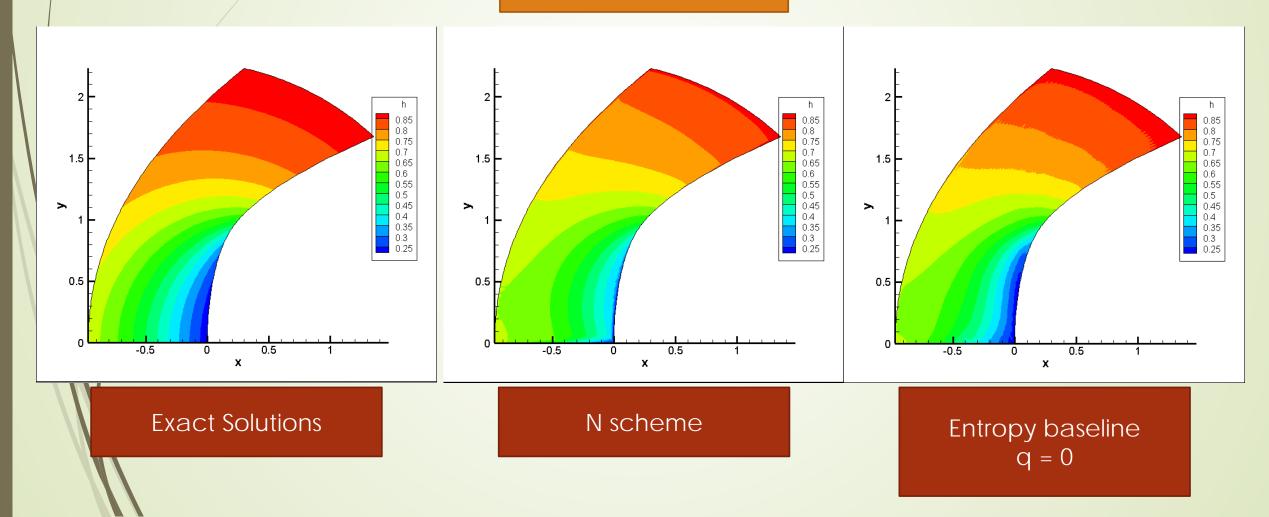


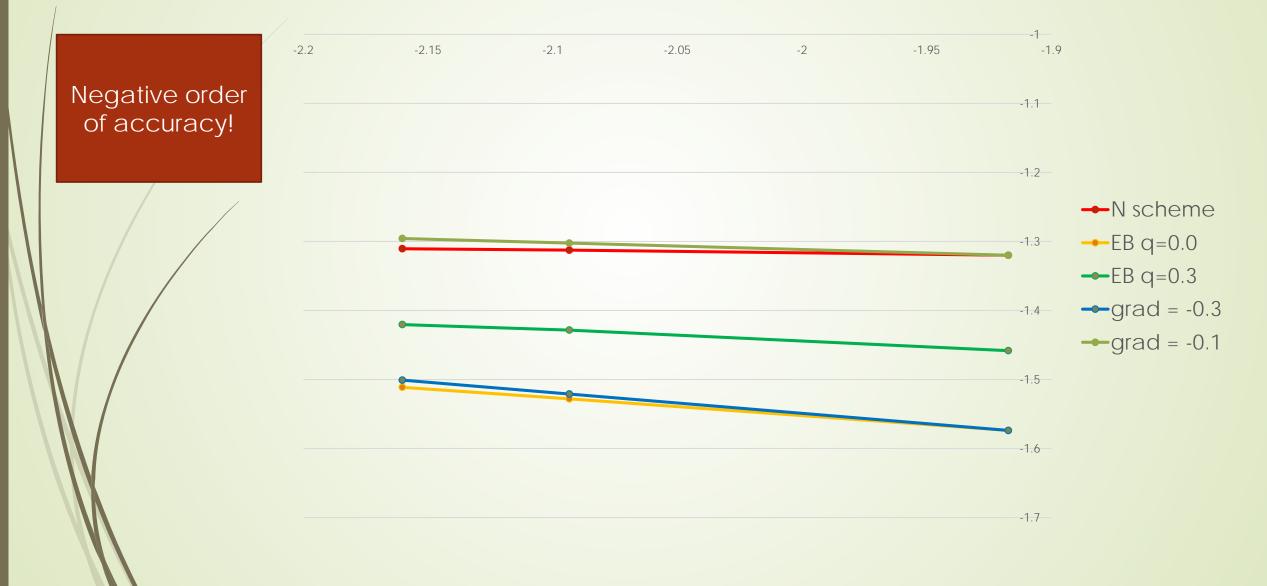
### Preliminary results

#### **Exact Solutions**

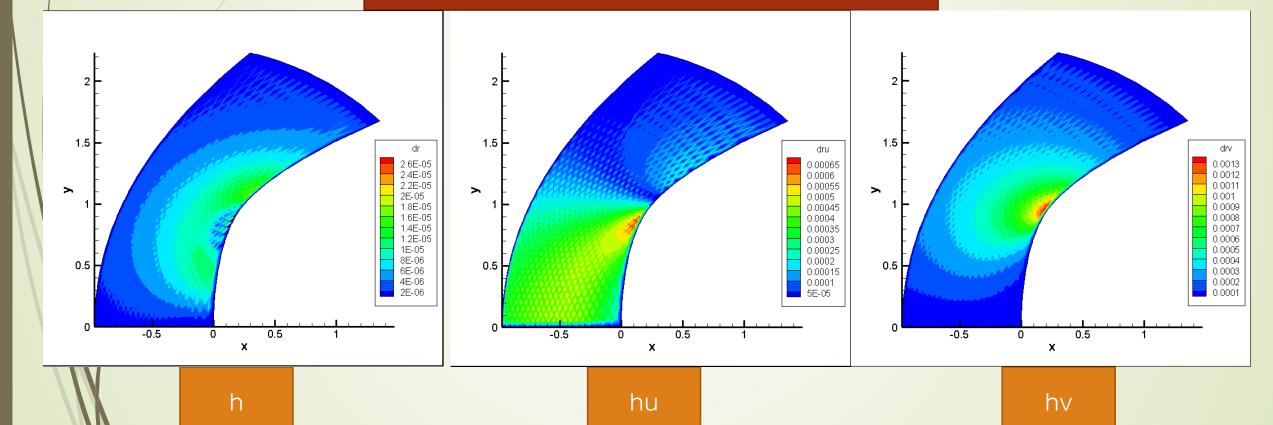


H (81 X 81 grid)





#### Error Contours at 1<sup>st</sup> iteration – N scheme



### Current proposals to resolve negative order of accuracy

- Double check the model with FV approach
- Inaccurate presentation of flow physics with Bernoulli equations?
- Incorporate hydraulic energy equation to SWE to preserve flow physics