Residual Distribution Schemes for Parabolic Equations

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- Focus is to incorporate viscous terms in the residual-distribution scheme
- Investigated in the context of 2-dimensional advection-diffusion problems

$$
u_t + au_x + bu_y = \nu(u_{xx} + u_{yy})
$$
\n(1)

 \bullet The semi-discrete form of the governing equation for node *i* is of the form:

$$
S_i \frac{du_i}{dt} + \sum_{T \in V_i} (\phi_{i, inv}^T + \phi_{i, vis}^T) = 0
$$
 (2)

• There are different approaches to evaluate the viscous component of the residual for each cell

• Begin with usual residual distribution approach for a triangle

$$
\phi_{\text{vis}}^T = -\iint_T \nu(u_{xx} + u_{yy}) dx dy \qquad (3)
$$

Using Gauss's theorem,

$$
\phi_{\mathsf{vis}}^{\mathcal{T}} = -\oint_{\partial \mathcal{T}} \nu \nabla u. d\mathbf{n} \tag{4}
$$

Since the solution is assumed to vary linearly over each triangle, the gradient for a triangle is constant

• This results in
$$
\phi_{\text{vis}}^T = 0
$$
 for the triangle

- Therefore we need a higher-order reconstruction and extension of the computational stencil are required to recover the ∇u on the boundary
- There are generally 2 alternatives
	- Nodal gradients using Least Squares approach
	- Edge based gradients using an arithmetic average approach with neighbouring triangle to recover the gradients at the midpoint
- However these approaches have the disadvantage of a larger stencil
- Another approach will be considered to keep the stencil compact

Viscous Component of Fluctuation Splitting

 \bullet Consider the median dual cell S_i

Figure: Median dual cell around node i

• Solve.

$$
\iint_{S_i} \nu(u_{xx} + u_{yy}) dx dy = \oint_{\partial S_i} \nu \nabla u \, d\mathbf{n}
$$
\n(5)\n
$$
\oint_{\partial S_i} \nu \nabla u \, d\mathbf{n} = \sum_{T \in V_i} \oint_{\partial S_i \cap T} \nu \nabla u \, d\mathbf{n}
$$
\n(6)

• Where the boundary $\partial S_i \cap T$ for each triangle T, is composed of 2 segments with outward scaled normals $\mathsf{n}^1_\mathrm{ext}, \mathsf{n}^2_\mathrm{ext}$ such that,

$$
\mathbf{n_{ext}^{1}} + \mathbf{n_{ext}^{2}} = -\frac{1}{2}\mathbf{n_i}
$$
 (7)

Viscous Component of Fluctuation Splitting

o Thus.

$$
\iint_{S_i} \nu(u_{xx} + u_{yy}) dx dy = \sum_{T \in V_i} \frac{\nu}{2} \nabla u . \mathbf{n_i}
$$
 (8)

• Where for a triangle T ,

$$
\nabla u = \frac{\sum_{p=1}^{3} u_p \mathbf{n_p}}{2S_T} \tag{9}
$$

- Note that the expression produced by the Galerkin method is identical
- Now the scheme is compact and a viscous component of the residual distribution for each triangle can be written equivalently as,

$$
\phi_{i,\text{vis}}^{\mathcal{T}} = \frac{1}{2} \mathbf{n}_i \cdot \nu \sum_{p=1}^3 \frac{u_p \mathbf{n}_p}{2S_{\mathcal{T}}}
$$
(10)

Viscous Component of Fluctuation Splitting

viscous term also has an influence on the time-step restriction, which becomes,

$$
CFL\frac{S_i}{\Delta t_i} = \sum_{T \in V_i} \frac{1}{\text{restri}_{inv}^i + \text{restri}_{vis}^i}
$$
(11)

Where $\mathit{restri}^i_{\mathit{vis}}$ is the diffusive time-step restriction (determined by positivity of the scheme), which is given by:

$$
restri_{vis}^i = \frac{\nu \|\mathbf{n_i}\|^2}{4S_T} \tag{12}
$$

- • Validation for the viscous residual using Galerkin approach for 201 nodal points in x and y directions and Dirichlet boundary conditions for the pure diffusion problem with $\nu = 1$.
- Where

$$
u_{\text{left}} = u_{\text{bottom}} = 0 \tag{13}
$$

$$
u_{right} = \sin(\pi y), u_{top} = -\sin(\pi x)
$$
 (14)

- Solving a pure diffusion problem for the domain $-1 \le x \le 0$ and $1 < y < 0$
- The analytic solution is:

$$
u(x,y) = \frac{1}{\sinh \pi} [\sinh(\pi(x+1))\sin(\pi y) + \sinh(\pi y)\sin(\pi(x+1))]
$$
 (15)

• Right running grid with skewness $= 0.3$

Figure: Close-up of the g[rid](#page-9-0)o \rightarrow 4 \oplus \rightarrow 4 \oplus

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Viscous Residual

• Numerical solution with skewness $= 0.3$

Figure: Exact solution Figure: Numerical solution

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Results

Comparison of the exact and numerical solution for 3 different y slices.

Figure: Comparison of the exact and numerical solution for $y = 0.25$

Figure: Comparison of the exact and numerical solution for $y = 0.50$

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Figure: Comparison of the exact and numerical solution for $y = 0.75$

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- • Solving for the domain $0 \le x \le 1$ and $1 \le y \le 0$ with $\nu = 0.1$ and $a = 7.0, b = 4.0$
- Using Dirichlet boundary conditions from the exact solution
- The exact solution is given as:

$$
u(x, y) = -\cos(\pi \eta) \exp[0.5\xi(1 - \sqrt{1 + 4\pi^2 \nu^2})/\nu]
$$
 (16)

Where

$$
\eta = bx - ay \tag{17}
$$

$$
\xi = ax + by \tag{18}
$$

• Inviscid residual computed using classic LDA scheme with grid skewness of 0.3

Figure: Exact solution for the linear advection-diffusion equation

Figure: Numerical solution with 101 nodes in x [an](#page-14-0)[d y](#page-16-0)[dir](#page-15-0)[e](#page-16-0)[ct](#page-13-0)[i](#page-14-0)[o](#page-20-0)[n](#page-21-0)

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Figure: Numerical solution with 201 nodes in x and y direction

Figure: Numerical solution with 301 nodes in x and y direction

Figure: Numerical solution with 401 nodes in x and y direction

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Comparison of the exact and numerical solution for 3 different y slices for 201 nodes in \times and \times directions.

Figure: Comparison of the exact and numerical solution for $y = 0.025$

Figure: Comparison of the exact and numerical [sol](#page-17-0)u[ti](#page-19-0)[o](#page-17-0)[n f](#page-18-0)[or](#page-19-0) $y = 0.15$ QQ

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 0.8

Comparison of the exact and numerical solution for 2 different x slices.

Figure: Comparison of the exact and numerical solution for $x = 0.015$

Figure: Comparison of the exact and numerical solution for $x = 0.95$

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Results

Figure: Comparison of the exact and numerical solution for $y = 0.75$ for 401 nodes

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Order of Accuracy

Table: Advection

Table: Advection-Diffusion

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Table: Order of Accuracy for Advection-Diffusion

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Table: L2-Norm Error Comparison

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Thank You

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