Alternative Approach to Prove Linear Preserving (LP) Property of RD Schemes

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2 [Flux-Difference Approach Proof](#page-4-0)

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- • The Linear Preserving (LP) property requires the scheme to preserve the exact steady state solution when the solution varies linearly in space for arbitrary grids
- A scheme that is LP will guarantee second-order accuracy for spatial discretization at steady state
- \bullet Recall the update for node i

$$
u_i^{n+1} = u_i^n + \frac{\Delta t}{S_i} \sum_{T \in V_i} \beta_i^T \phi^T
$$
 (1)

where $\phi^\mathcal{T}$ is the total residual and β_i is the distribution coefficient

Basic Proof 2

- At steady state, is we substitute the exact solution at the nodes, we get $\phi^{\mathcal{T}}=0$
- If we assume that β_i is bounded as $\phi^\mathcal{T}$ approaches zero, the from Eq. [1](#page-2-1) we get

$$
u_i^{n+1} = u_i^n \tag{2}
$$

- Thus, we preserve the exact solution which proves LP
- Conversely, a change in u_i is only possible if $\beta_i^{\mathcal T} \phi^{\mathcal T}$ is not zero
- This implies that β_i is not bounded as $\phi^\mathcal{T}$ approaches zero
- Thus, boundedness of the distribution coefficient is a sufficient condition for LP
- For example, LDA distribution coefficients defined as $\beta_i = \frac{k_i^+}{\sum k^+}$ is always between 0 and 1 (thus bounded)

• Recall, the flux difference signal to node i is

$$
\phi_i^{FD} = \frac{1}{2}(\vec{f}_i - \vec{f}^*) \cdot \vec{n_i} + A.T
$$
 (3)

- o where A.T are the artificial terms
- Naturally, to prove LP, we would need to cast this approach in the classic RD form which is

$$
\phi_i^{Flux-Diff} = \beta_i^{Flux-Diff} \phi^{\mathsf{T}}
$$
 (4)

where

$$
\beta_i^{Flux-Diff} = \frac{\frac{1}{2}(\vec{f}_i - \vec{f}^*) \cdot \vec{n_i} + A.\mathsf{T}}{\phi^{\mathsf{T}}} = \frac{(2ui - uj - uk)k_i + A.\mathsf{T}}{\phi^{\mathsf{T}}} \quad (5)
$$

- From Eq. [5,](#page-4-1) as $\phi^\mathcal{T}$ approaches zero, $\beta_i^\mathit{Flux-Diff}$ i ^{Flux-Diff} is unbounded
- **•** From the reasoning above, the flux-difference approach is not LP
- However, from numerical experiments we know that the flux-difference approach is 2nd order accurate and should be LP
- We first conclude that we can't interpret the flux-difference approach as done in Eq. [4](#page-4-2)
- To prove LP for flux-difference approach, we need an alternate approach

 \bullet At steady state, the sum of signals to node *i* from all the triangular elements that share the node is,

$$
\sum_{T,i\in\mathcal{T}}\phi_i^{\mathcal{T}}=0.\tag{6}
$$

Introducing a smooth function $\Phi \in \mathcal{C}^1$ and take the product with Eq. [6](#page-6-0) as well as taking the summation over all nodes, the following relation is obtained.

$$
\sum_{i} \Phi_{i} \cdot \sum_{T, i \in T} \phi_{i}^{T} = 0 \tag{7}
$$

• Recall for a triangular element T ,

$$
\sum_{i \in T} \phi_i^T = \iint_T \vec{\nabla} \cdot \vec{F}^h dA. \tag{8}
$$

• Introduce for a triangular element T with vertices/nodes i, j, k a function $\Phi^{\mathcal{T}}$ defined as $\Phi^{\mathcal{T}} = \frac{\Phi_i + \Phi_j + \Phi_k}{3}$ $\frac{\nu_j+\nu_k}{3}$ and take the product with Eq. [8](#page-6-2) which results in,

$$
\sum_{T} \Phi^{T} \sum_{i \in T} \phi_{i}^{T} = \sum_{T} \Phi^{T} \iint_{T} \vec{\nabla} \cdot \vec{F}^{h} dA.
$$
 (9)

Now subtracting the terms in Eq. [7](#page-6-3) from Eq. [9,](#page-7-1) we obtain

$$
\sum_{T} \Phi^{T} \iint_{T} \vec{\nabla} \cdot \vec{F}^{h} dA + \sum_{T} \sum_{i \in T} (\Phi_{i} - \Phi^{T}) \cdot \phi_{i}^{T} = 0 \qquad (10)
$$

To obtain second-order accuracy at steady state, the second terms in Eq. [10](#page-7-2) must be of order h^2 . Which is,

$$
\sum_{i} (\Phi_{i} - \Phi^{\mathcal{T}}) \cdot (\sum_{\mathcal{T}, i \in \mathcal{T}} \phi_{i}^{\mathcal{T}}) = O(h^{2})
$$
 (11)

This is true when the signal $\phi_i^{\mathcal{T}}$ is $O(h^3)$ since for a bounded domain the number of node[s](#page-4-0) (\sum_i) (\sum_i) [i](#page-3-0)s $O(h^{-2})$ $O(h^{-2})$ $O(h^{-2})$ $O(h^{-2})$ and $(\Phi_{\hat{t}}-\Phi_{\widehat{\epsilon}}^{\mathcal{T}})$ is $O(\underline{h})$ [.](#page-10-0)

As an example, signal for classic LP(bounded distribution coefficients β_i) RD is of the form,

$$
\phi_i^{\mathcal{T},\mathcal{LDA}} = \beta_i \phi^{\mathcal{T}}.\tag{12}
$$

- Recall that $\phi^{\cal T}=\oint\vec{F}\cdot\hat{n}dl$ and approximating the flux with trapezoidal rule which is $O(\mathit{h}^{2})$ and taking the product with the scaled normals $(O(h))$ gives the $\phi^{\mathcal{T}} = O(h^3)$
- Thus, with bounded coefficients, the signal will be $O(h^3)$.
- For the flux-difference approach, the baseline approach ignoring \vec{f}^* where the residual over an element is evaluated using a trapezoidal rule and is of the form,

$$
\phi^{\mathcal{T}} = \frac{1}{2}\vec{f}_i \cdot \vec{n}_i + \frac{1}{2}\vec{f}_j \cdot \vec{n}_j + \frac{1}{2}\vec{f}_k \cdot \vec{n}_k. \tag{13}
$$

From Eq. [13](#page-8-1) above, it is clear that residual, $\phi^{\, \mathcal{T}} = O(h^3).$ The signal to the nodes will be,

$$
\phi_{i,j,k} = \frac{1}{2} \vec{f}_{i,j,k} \cdot \vec{n}_{i,j,k} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i\}} \mathcal{A} \cdot \sum_{\text{all } i \in \mathbb{N} \setminus \{i
$$

- Since each signal is a portion of the total residual, $\phi^\mathcal{T}$, thus each signal will also be $O(h^3)$
- The basic flux-difference approach is second-order accurate
- With the choice of \vec{f}^* being the arithmetic average, the overall second-order accuracy is also preserved
- This because \vec{f}^* is defined as arithmetic average of the average of the fluxes along each edge which is,

$$
\vec{f}^* = \frac{1}{3} \left(\frac{1}{2} (\vec{f}_i + \vec{f}_j) + \frac{1}{2} (\vec{f}_j + \vec{f}_k) + \frac{1}{2} (\vec{f}_k + \vec{f}_j) \right)
$$

= $\frac{1}{3} (\vec{f}_i + \vec{f}_j + \vec{f}_k)$ (15)

- From Eq. [15](#page-9-1) above, since \vec{f}^* is constructed as an arithmetic average of using trapezoidal rule (which is 0($h^2)$) along each edge, \vec{f}^* will be $O(h^2)$
- As a result, \vec{f}^* preserves the order of accuracy the full flux-difference approach is LP on arbitrary grids. Ω

• The viscous signal can be written as,

$$
\phi_{i,\text{vis}}^T = \frac{\nu}{4A_T} \vec{n}_i \cdot \vec{n}_j (u_j - u_i) + \frac{\nu}{4A_T} \vec{n}_i \cdot \vec{n}_k (u_i - u_k). \qquad (16)
$$

• Now, writing the artificial signals in the same form by setting $\beta = 0$ we obtain,

$$
\phi_i^{\text{art}} = \alpha (u_j - u_i) + \gamma (u_i - u_k). \tag{17}
$$

- Comparing Eq. [16](#page-10-1) and Eq. [17,](#page-10-2) we conclude that the artificial signals are a Galerkin type discretization since the terms multiplying $(u_i - u_i)$ (which are the dot product of the scaled normals) in Eq. [16](#page-10-1) are of $O(h^2)$ while α, γ are of $O(h^2)$
- This similar form leads to the conclusion that since the Galerkin discretization is second-order, the artificial signal too will be second-order accurate

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